

NUMERICAL EXPERIMENTS OF SINGLE MODE GYROTRON EQUATIONS¹

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The present work continues our recent investigations of the stationary problem of the single mode gyrotron equation [1] using the implicit finite difference schemes and the method of lines. We consider two versions of the gyrotron equation. The amplitude $f(t, x)$ of the high frequency field in the gyrotron resonator and the transverse orbital momentum $p(T, x, \theta_0)$ of electrons can be described by the following system of two complex differential equations (new version):

$$\begin{cases} \frac{\partial p}{\partial x} + i(\Delta + |p|^2 - 1 - g_b)p = if(t, x) \\ \frac{\partial^2 f}{\partial x^2} - i(1 + \delta_\omega)\frac{\partial f}{\partial t} + (1 + 0.5(\delta_\omega + g_c))g_d f = (1 + \delta_\omega)(1 + g_c)^2 I \langle p \rangle, \end{cases} \quad (1)$$

where $i = \sqrt{-1}$, $x \in [0, L]$ and $t \geq 0$ are the axial and time coordinates, L - the length of the interaction space, $\Delta, \delta_\omega, \theta_0 \in [0, 2\pi]$ - the real constants, I - the current, $g_b(x), g_c(x), g_d(x)$ - given real functions and $\langle p \rangle = \frac{1}{2\pi} \int_0^{2\pi} p d\theta_0$ averaged value of p . The system (1) is supplemented by the initial conditions $p(t, 0, \theta_0) = \exp(i\theta_0)$, $f(0, x) = f_0(x)$, and by the boundary conditions in the gyrotron cavity

$$f(t, 0) = 0, \quad \frac{\partial f(t, L)}{\partial x} = -i\gamma f(t, L),$$

where $f_0(x)$ is the given complex function, γ is a positive parameter.

REFERENCES

- [1] J.Cepītis, H.Kalis and A. Reinfelds. Numerical investigations of single mode gyrotron equations. *Mathematical Modelling and Analysis*, **14** 2:169-178, 2009.

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